



# Glass Transition Distribution in Miscible Polymer Blends: from Calorimetry to Rheology

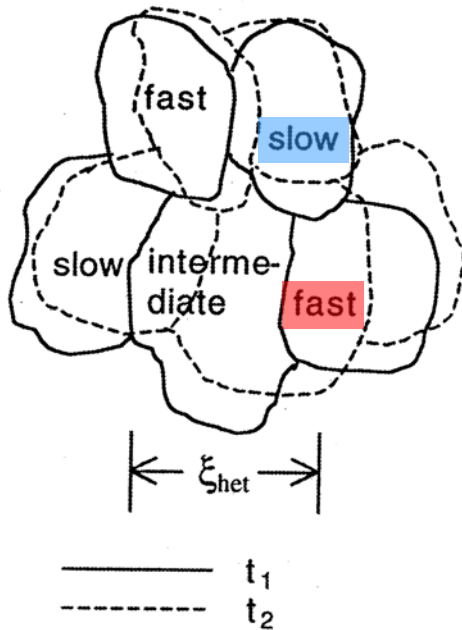
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E.S.P.C.I. ParisTech  
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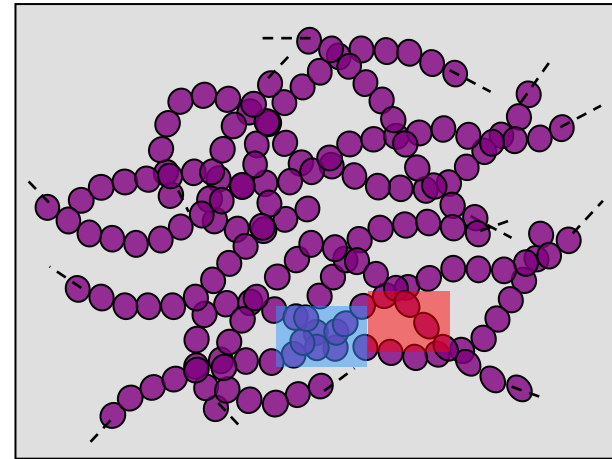
Macromolecules in Constrained Environments  
March 24-29<sup>th</sup>, 2013

# Classic picture of glasses: “dynamic heterogeneity”

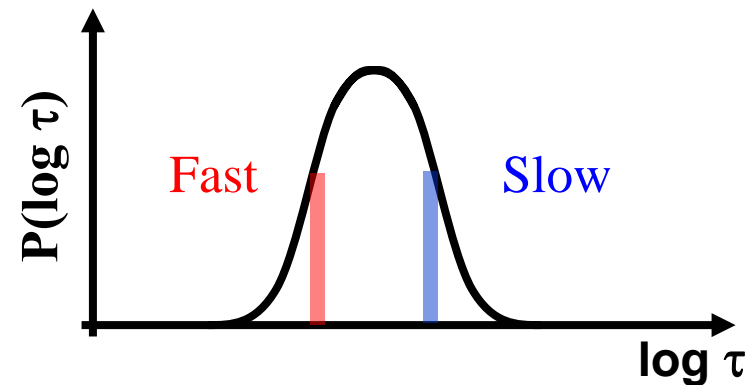
[Ediger, Review, 2000]



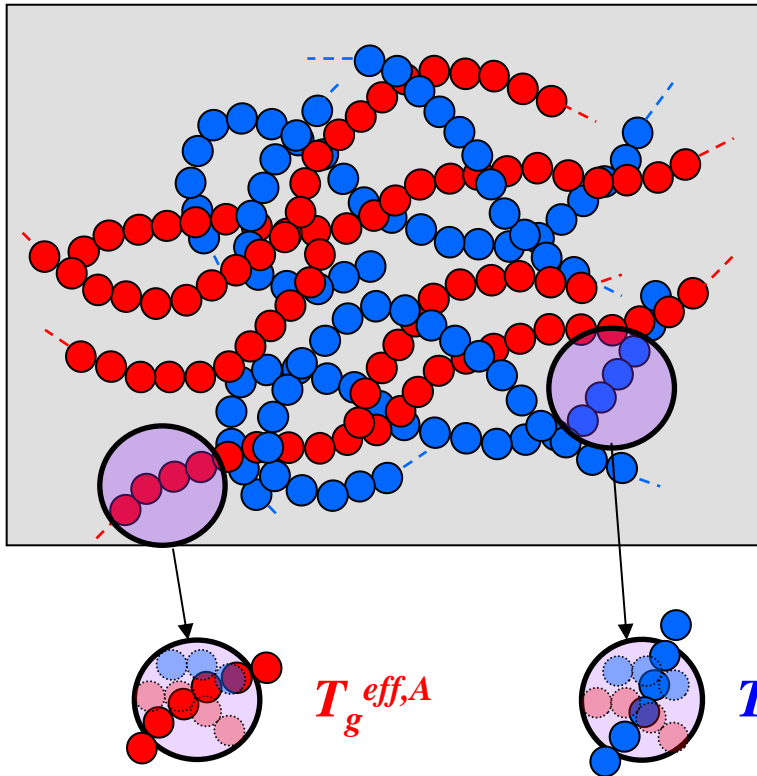
« Schematic illustration of regions of spatially heterogeneous dynamics near  $T_g$ . These regions are on the order of  $\xi_{het}$  in dimension (typically a few nanometers) and evolve in time. »



A glassy polymer



# Classic picture of miscible polymer blends: dynamic asymmetry



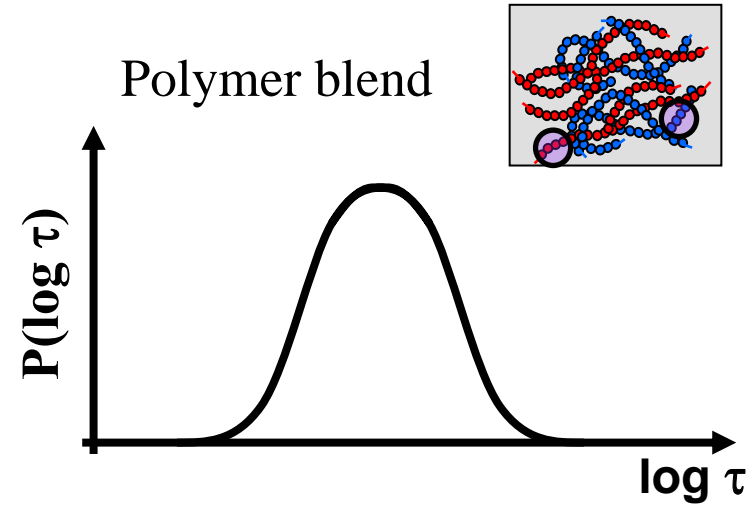
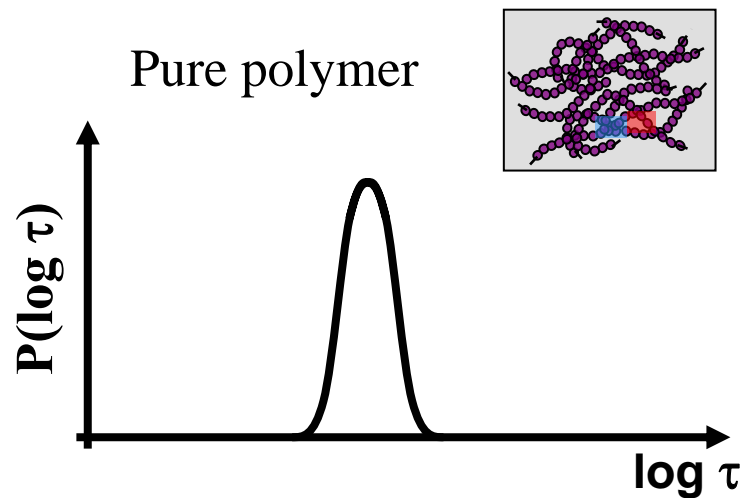
Fast component, low  $T_g$

Slow component, high  $T_g$

Great  $\Delta T_g \rightarrow$  dynamic asymmetry

Self-concentration (Lodge & McLeish, Colby, etc...)

# Miscible polymer blends: dynamic asymmetry



→ Glass Transition Distribution??

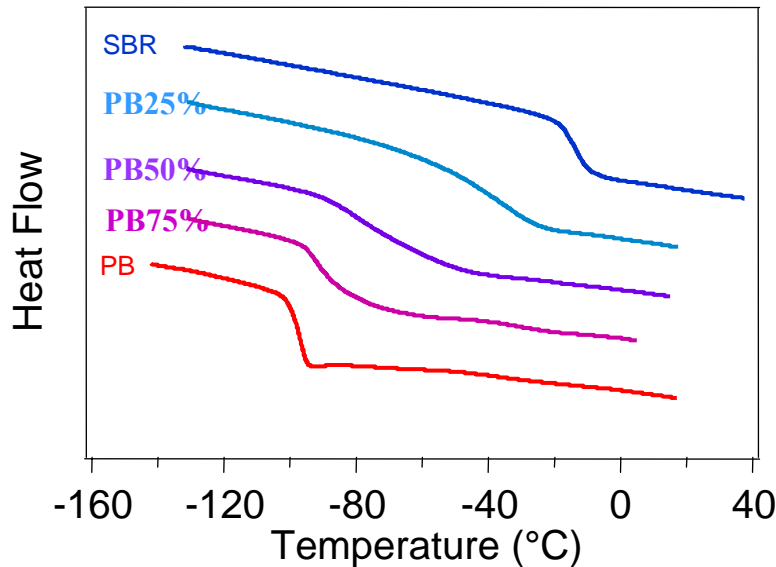
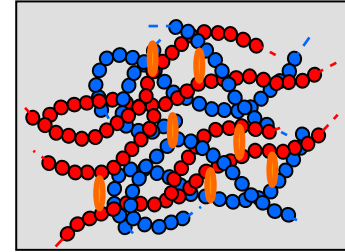
- Calorimetry
- Physical aging
- Rheology

# Outline

- From calorimetry to  $P(T_g)$  of blends
- Physical aging of blends
- Rheology prediction
- Conclusion

# Our blends PB/SBR

- PB: Polybutadiene,  $T_g \sim -100^\circ\text{C}$ ,  $l_k \sim 10\text{\AA}$
  - SBR: random Styrene-Butadiene copolymer,  $T_g \sim -15^\circ\text{C}$ ,  $l_k \sim 12\text{\AA}$
- All blends are mixed in common solvent and crosslinked.



PI/PVE blends

[Lodge & McLeish, 2000]

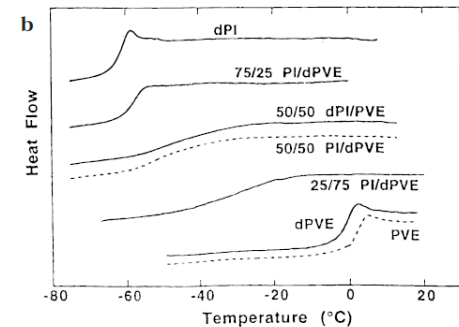


Figure 5. (a) Effective glass transition temperatures and (b) dsc traces for PI/PVME blends; data from Chung et al.<sup>13</sup>

→ A single and very broadened  $T_g$  for our blends.

# From Calorimetry to $P(Tg)$

$P_{deriv}(Tg) \sim$  Temperature derivative of the heat flow.

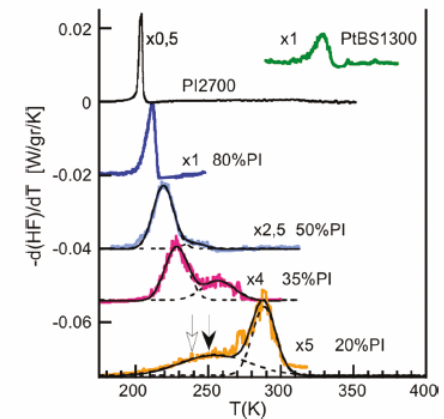
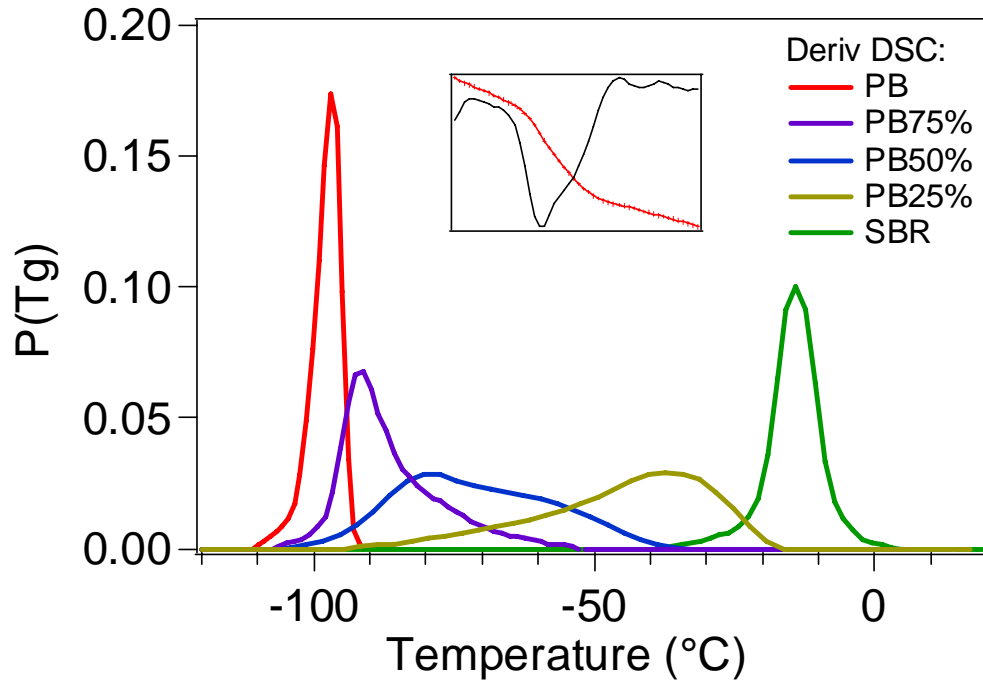


Figure 1. Lower panel: Derivative of the heat flow with respect to T for PI2700 and PtBS1300 homopolymers and their blends at different compositions (color on line). The intensity of the curves was multiplied

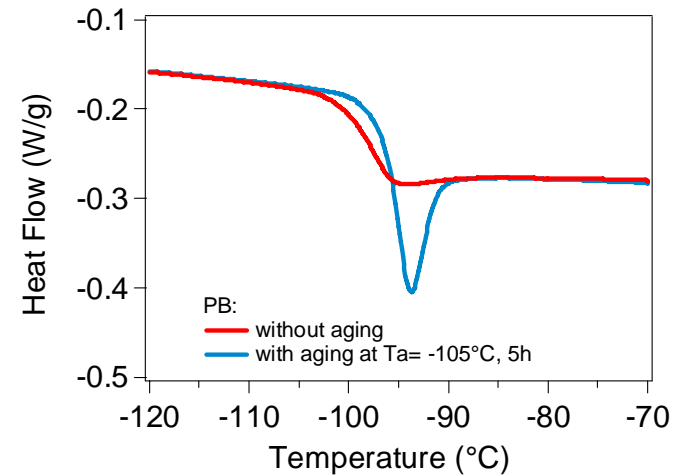
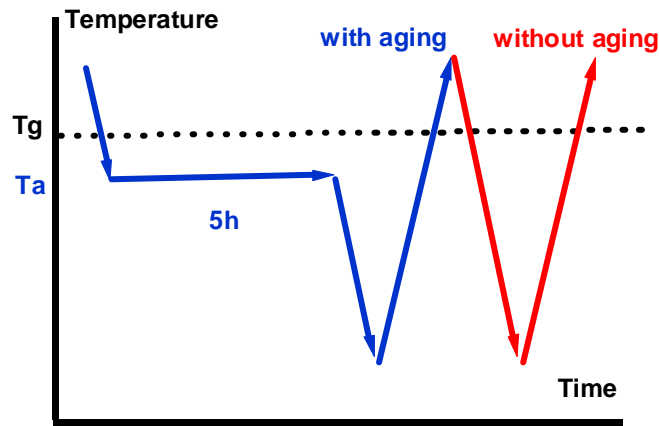
[Arrese-Igor, Alegria, and Colmenero, 2010]

➔ A single and very broadened  $Tg$  for our blends.

➔ What about physical aging ?

# Physical aging of homopolymers

## First observation DSC:

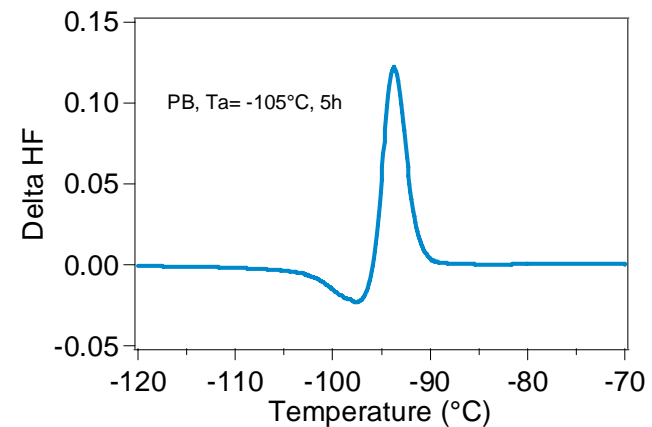


Annealing temperature:  $T_a$

Annealing time  $\tau_a = 5\text{h}$

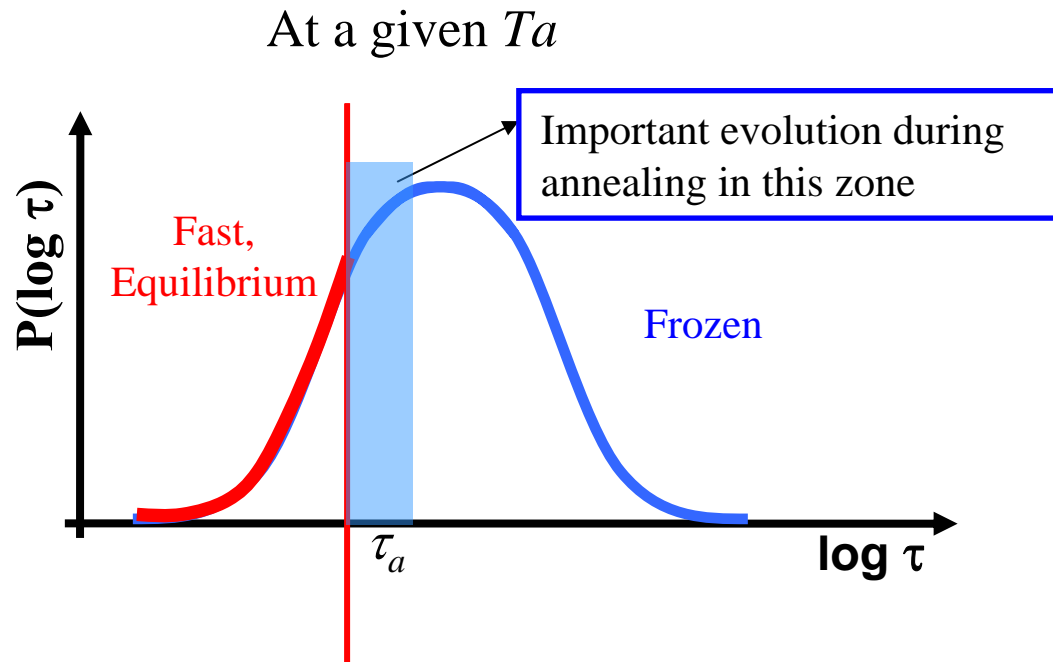
→ An overshoot peak!

→ Why?





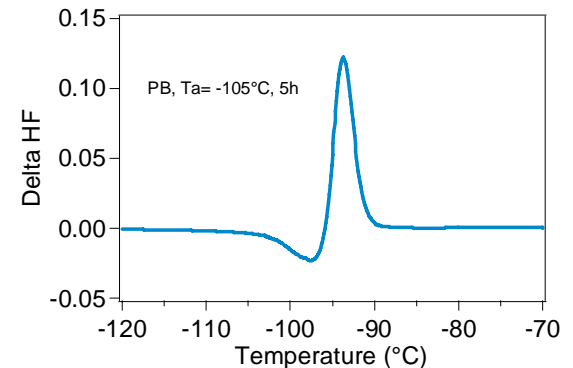
# What happened during annealing?



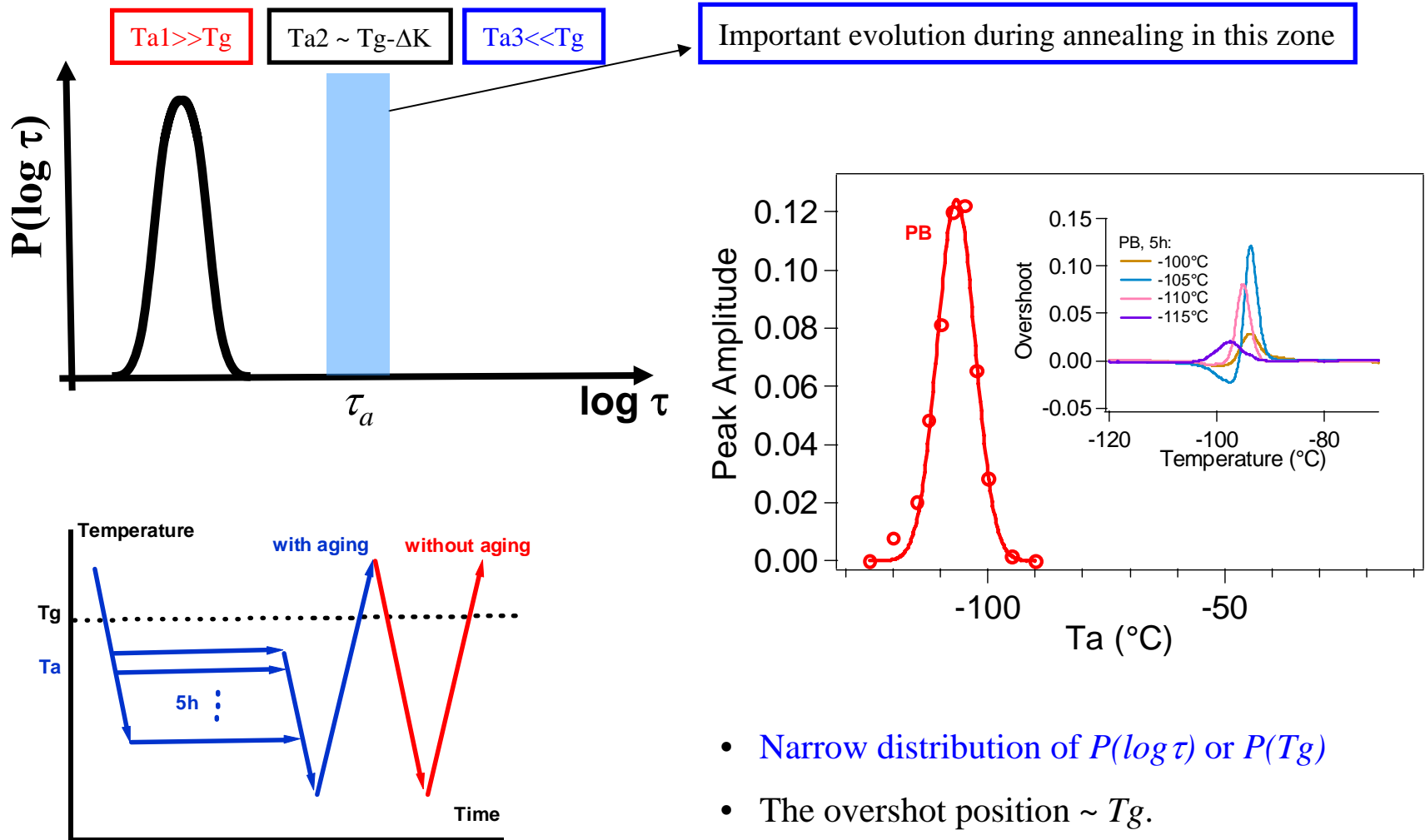
Evolving towards **longer** relaxation times.  
Minima energy landscape.

➔ Origin of the overshoot peak.

➔ If we change  $T_a$ ??



# Physical aging of homopolymers

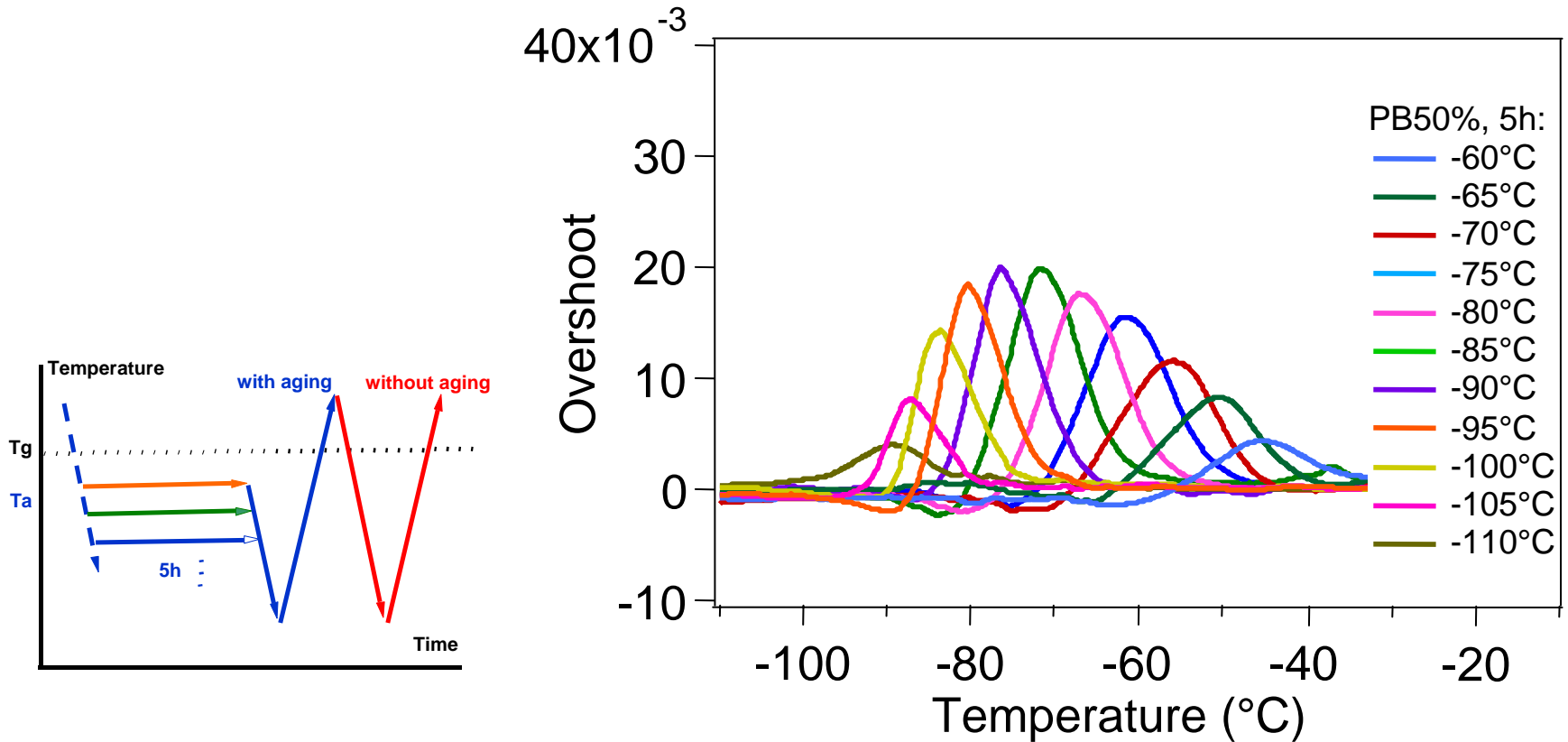


Important evolution during annealing in this zone

- Narrow distribution of  $P(\log \tau)$  or  $P(Tg)$
- The overshoot position  $\sim Tg$ .

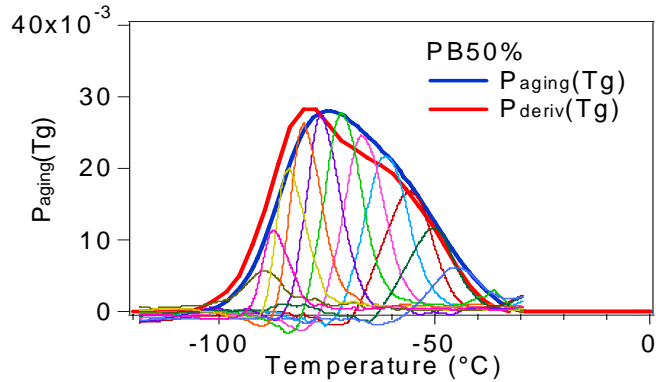
➔ For blends??

# Physical aging of blends

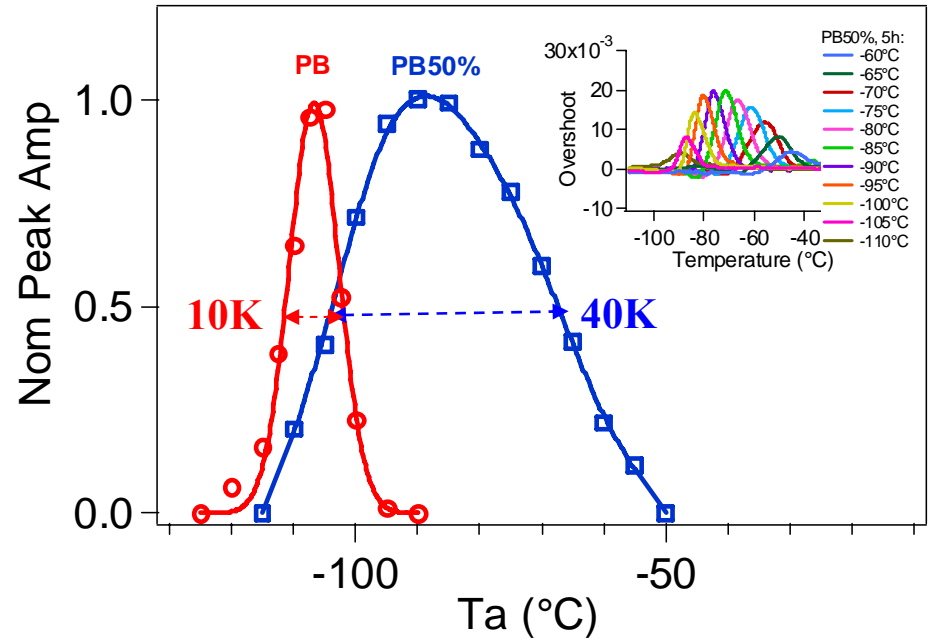
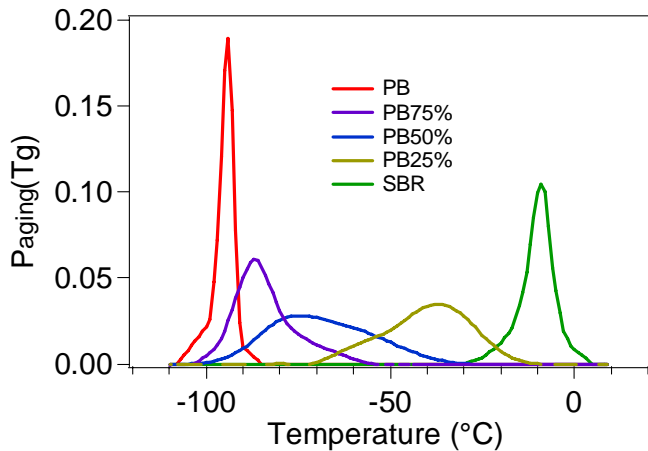


A broad distribution of overshoot peaks → broad  $P(T_g)$  ??

# Physical aging of blends

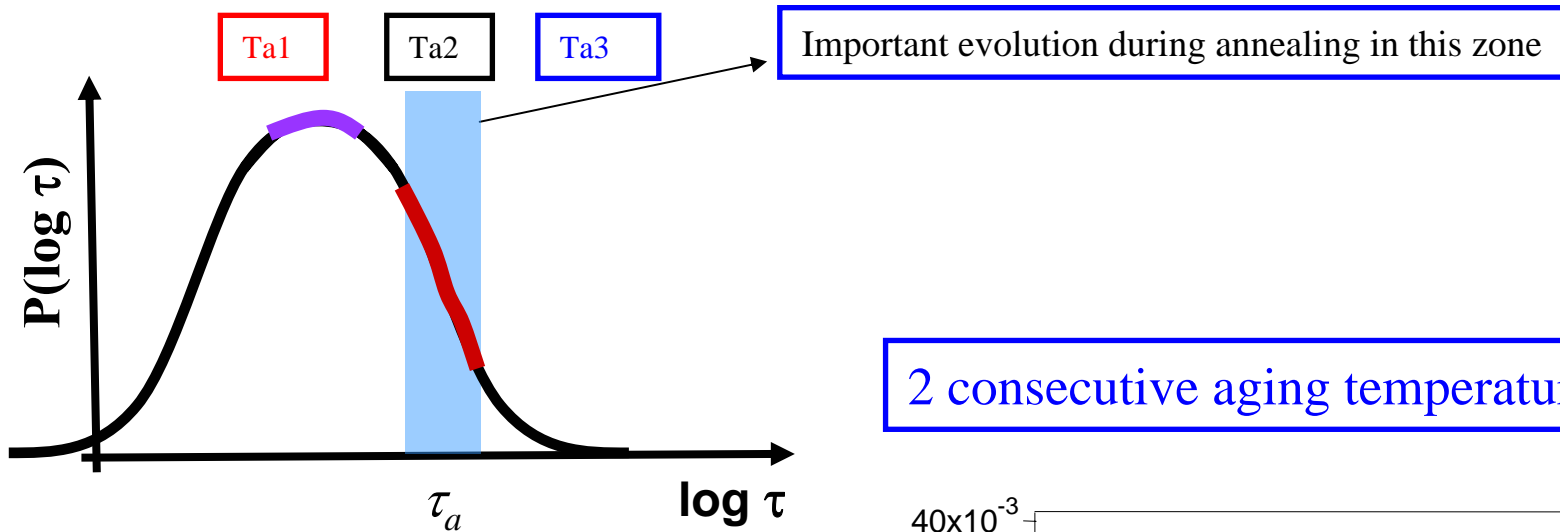


The envelop  $\sim P_{aging}(Tg) \sim P_{deriv}(Tg)$

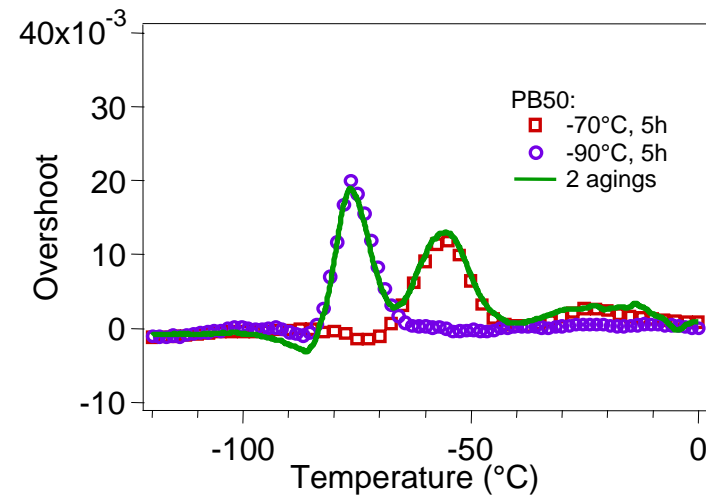
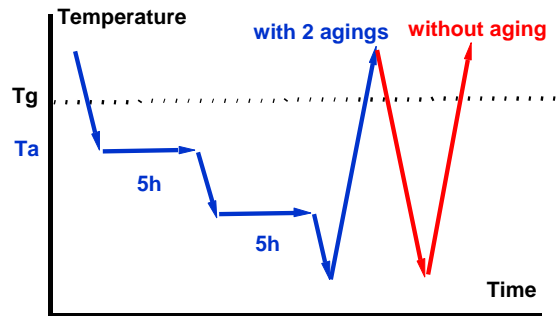


- Broad distribution of  $P(Tg)$
- Different domains are selected??

# Physical aging of blends

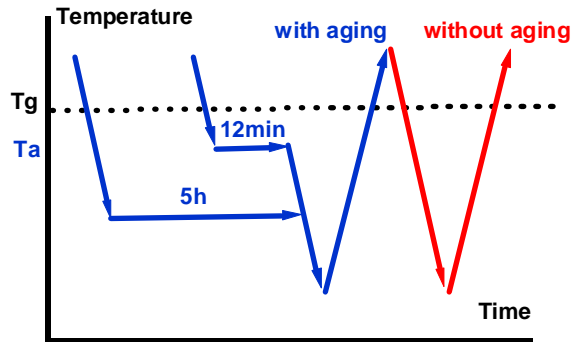


2 consecutive aging temperatures??



→ Each aging overshoot is controlled independently by a corresponding  $T_a$

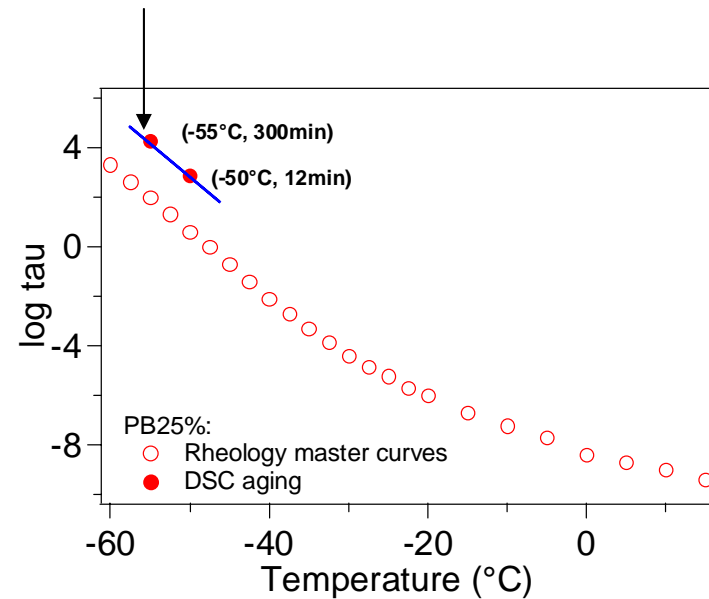
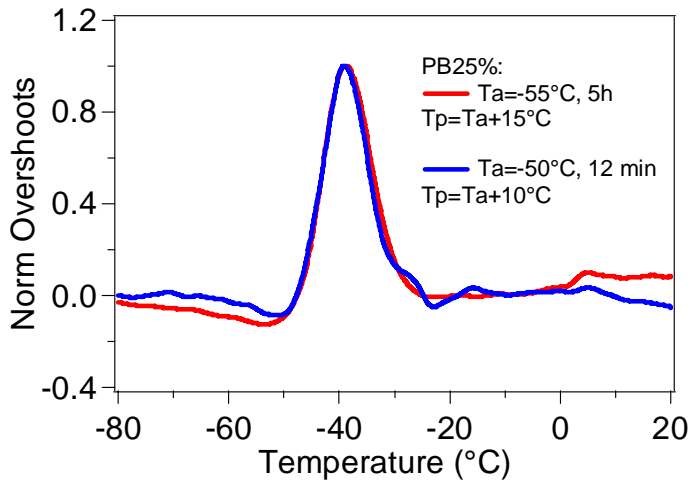
# $T_a - \tau_a$ superposition



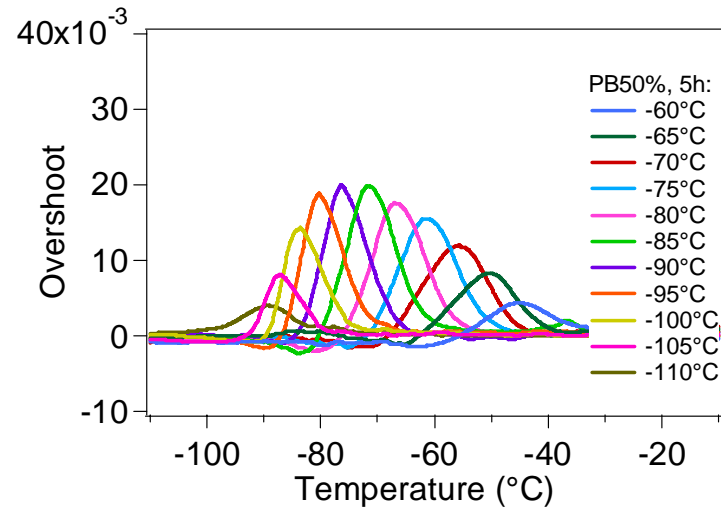
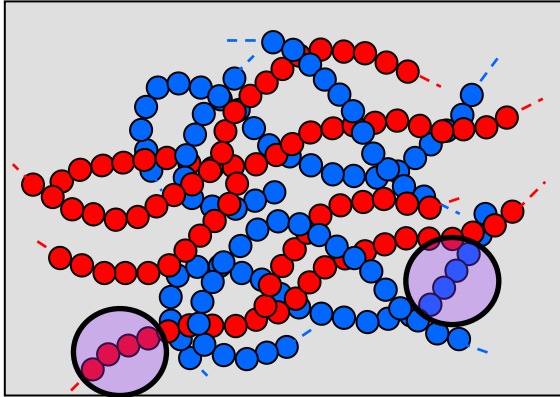
$T_a - \tau_a$  superposition...

$(-55^\circ\text{C}, 5\text{h}) \sim (-50^\circ\text{C}, 12\text{min})$

Same slope  $\rightarrow$  similar time temperature superposition law as rheology (WLF)



# Physical aging measurement

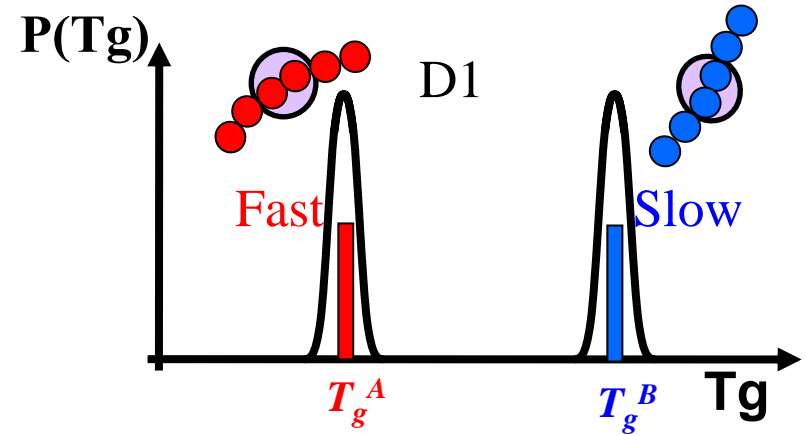
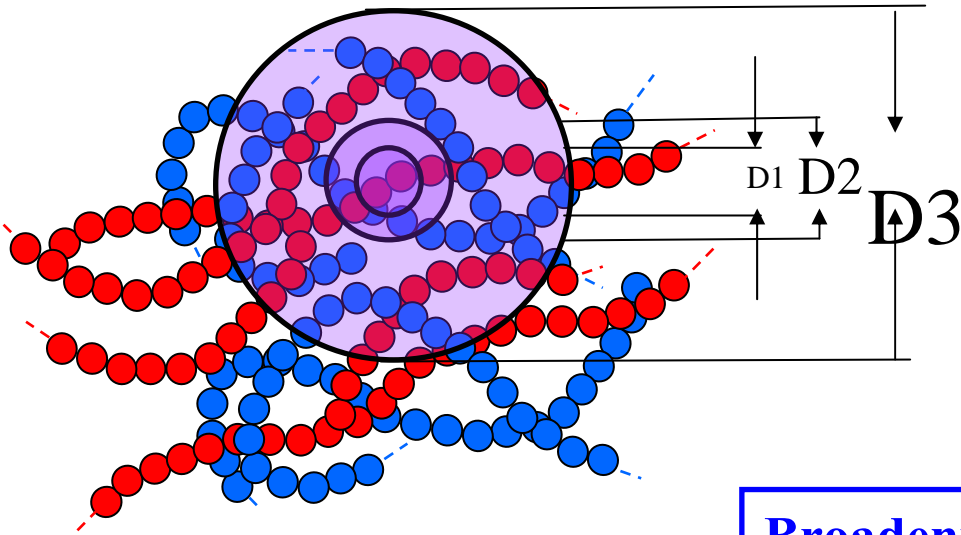


## Conclusion 1:

- DSC measurement with physical aging at different  $T_a$  exhibits a peak that reveals the independent contribution of different domain of  $T_g$ .
- The  $\tau_a$ - $T_a$  superposition seems to be valid, and is similar to rheology t-T superposition.
- The envelop  $\sim P_{aging}(Tg) \sim P_{deriv}(Tg)$

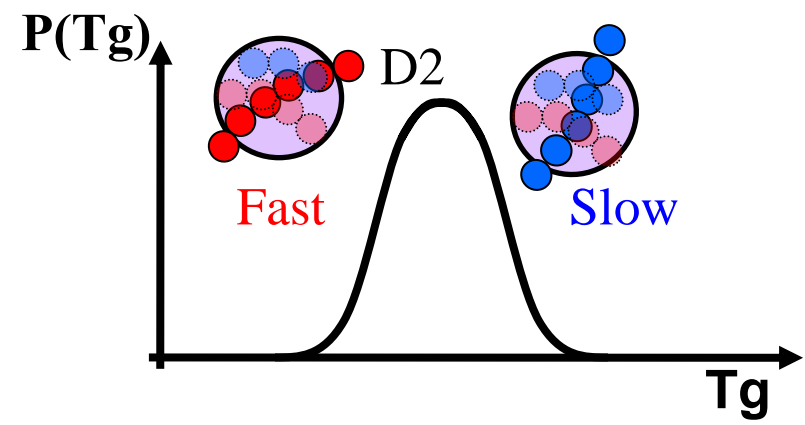
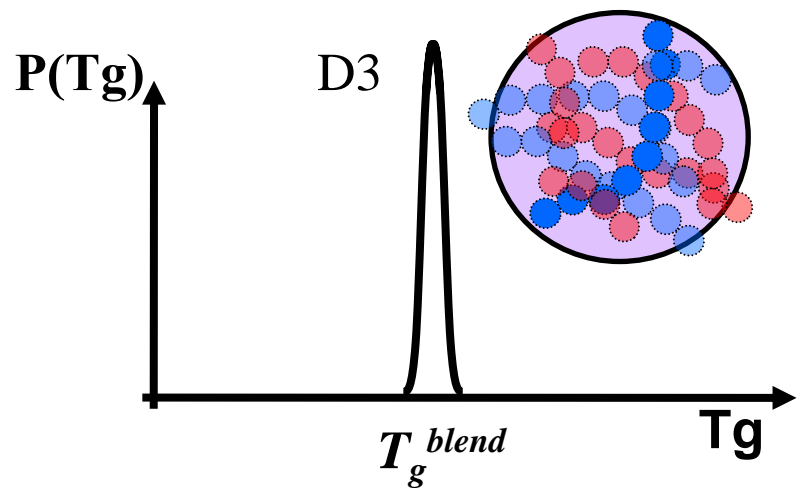
➔ The size scale of the glass transition domain??

# The size scale $D$



Broadening  $\propto 1/D$

$$\frac{1}{T_g(\varphi)} = \frac{\varphi}{T_g^{PB}} + \frac{1-\varphi}{T_g^{SBR}}$$





# The size scale $D$

Estimation of a characteristic size scale  $D$  using Lodge & McLeish + Shenogin & Colby method, we got  $\xi \sim 1\text{nm}$  (a typical value).

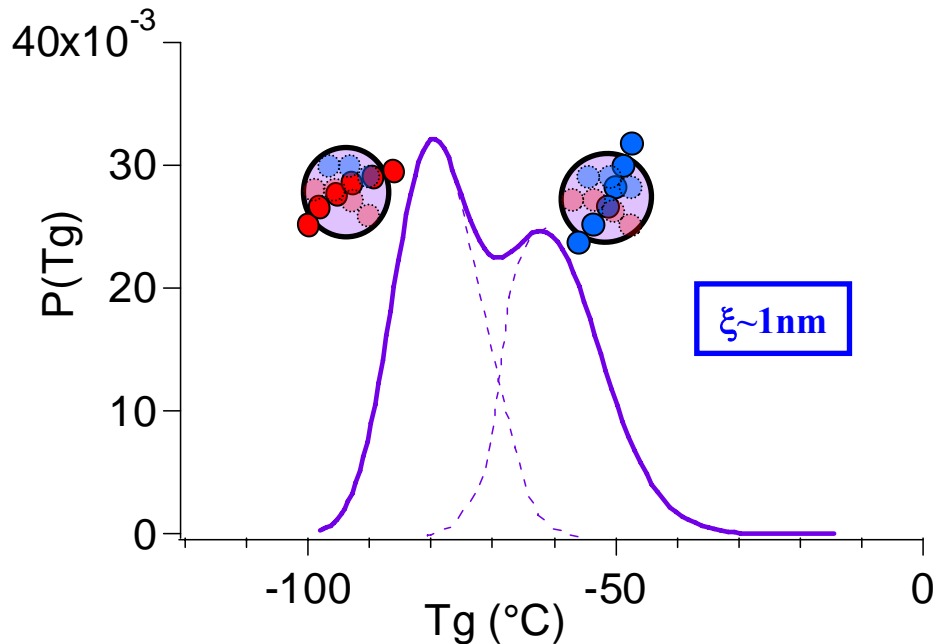


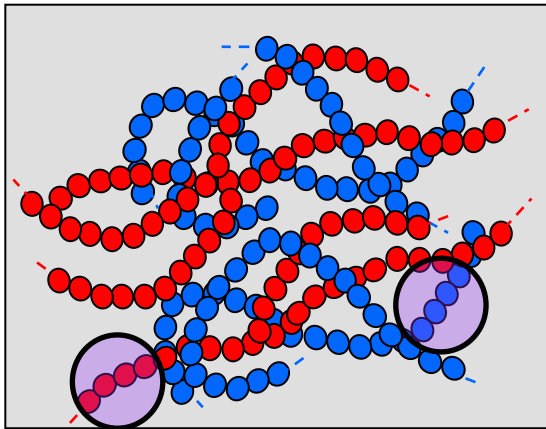
Table 2. Average Values of Correlation Radius and Respective Values of Self-Concentration, Found from the Fits of Eq 20 to the Dielectric Relaxation Data<sup>a</sup>

components	$R_c$ (Å)	$\phi_{self}$	$T_g$ (K)
PI	4	0.78	210
PVE	8	0.29	273
PBO	4.3	0.51	200
PVE-89	8	0.29	269
PVME	5.4	0.57	249
PS	$\sim 25$	0.053	379
P2CIS	16	0.13	402

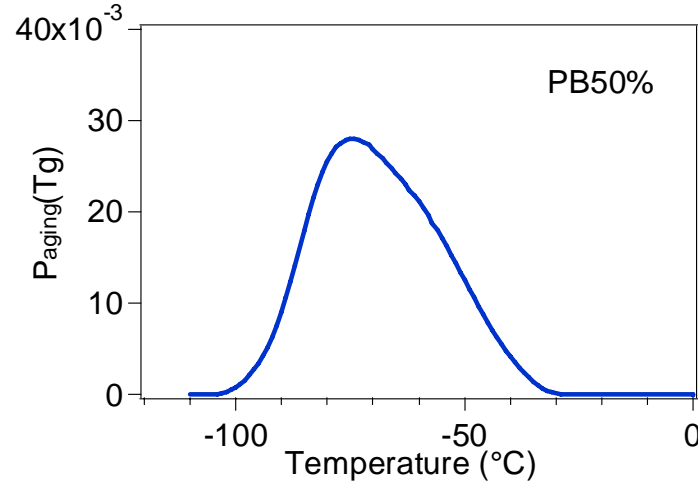
[Shenogin et al., 2007]

**Conclusion 2:** Our blend system can be considered as an ensemble of domains with different glass transition temperature, and the size scale is  $\sim 1\text{nm}$ .

# Rheology prediction



DSC

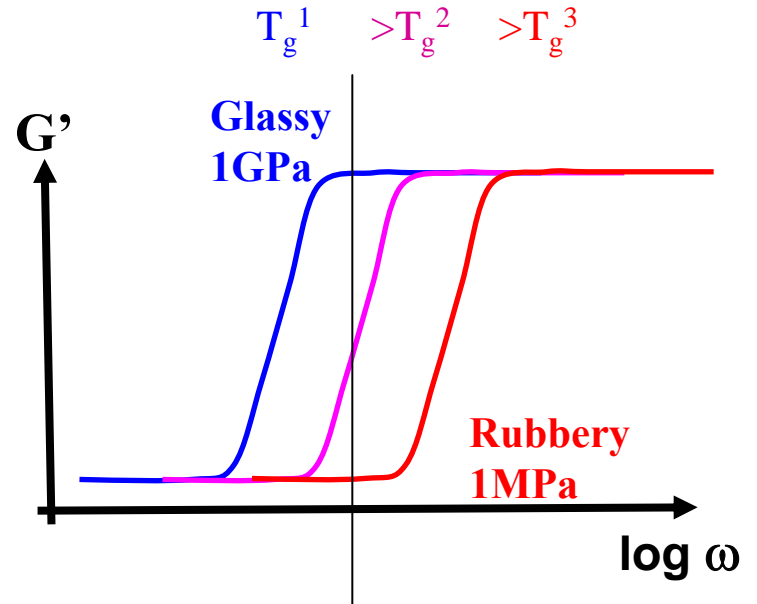
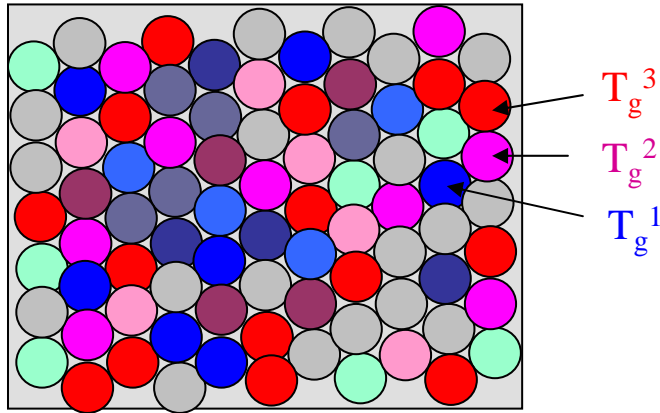


Rheology?

→ Is the rheology property controlled by the same size scale, and the same Tg distribution??

# Rheology prediction

$P(\log \tau)$  or  $P(Tg)$



Viscoelastic spectra of a local domain of  $G^*(Tg)$ :

WLF law + modified Havriliak Negami function

$$G^* = G_{glass} - \frac{G_{glass} - G_{rub}}{\left(1 + (j\omega\tau_{HN})^\alpha\right)^\beta}$$

All the parameters are interpolated from pure polymers at various local  $Tg$

$G'_{macro} = \text{average of } G'_i$

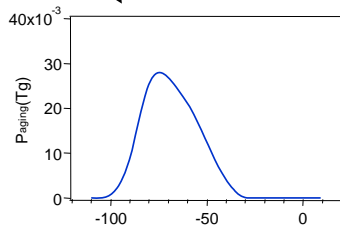
➔ HOW??

# Parallel or in-series?



$$G_{//}^* = \int G^*(T_g) P(T_g) dT_g$$

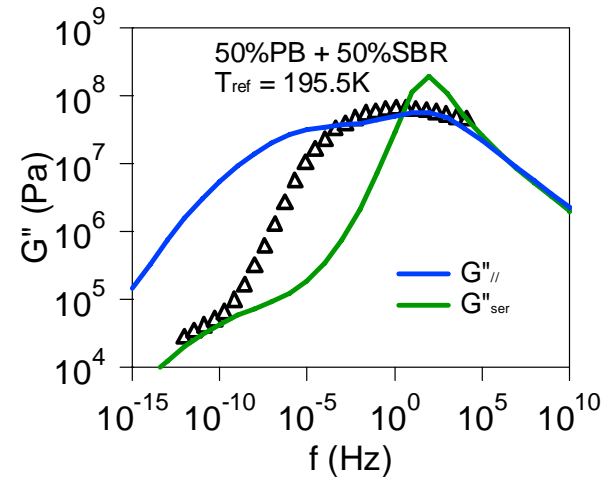
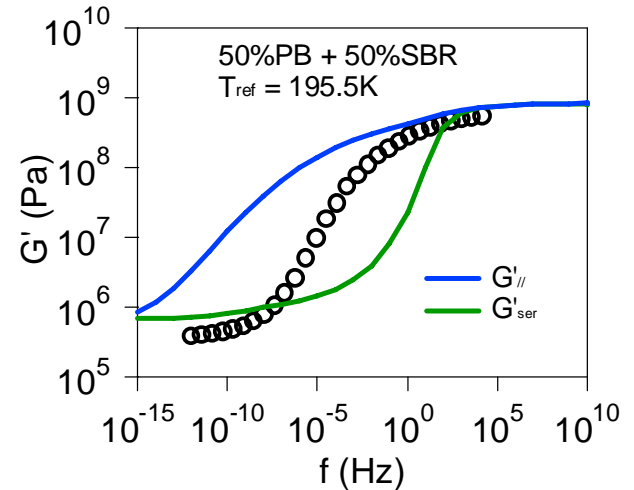
$$\frac{1}{G_{ser}^*} = \int \frac{P(T_g)}{G^*(T_g)} dT_g$$



*distributions from  $P_{aging}(Tg)$*

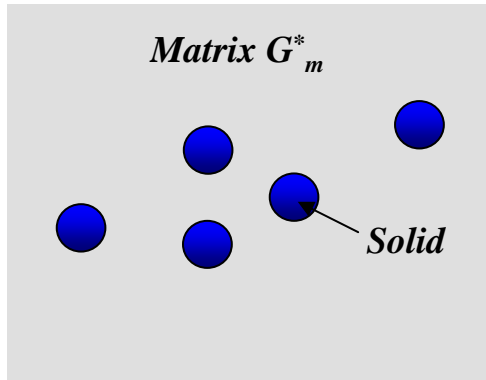
➔ Neither the parallel nor the series averaging is able to give a correct description.

➔ Need for a more refined model.



Rheology data: master curves

# Rheology prediction



Dilute solid particle-filled system:

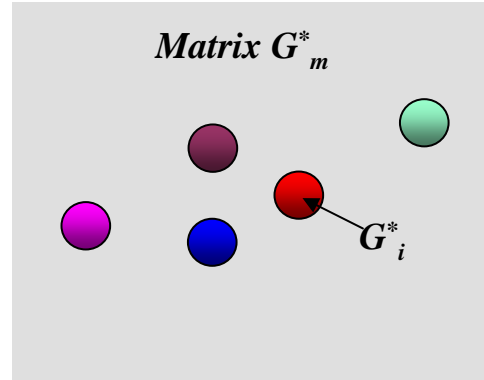
$$G^* = G_m^* \left( 1 + \frac{5}{2} \Phi \right)$$

Volume fraction

Einstein equation.

[Einstein, 1911]

[Smallwood, 1944]



Dilute visco-elastic inclusions:

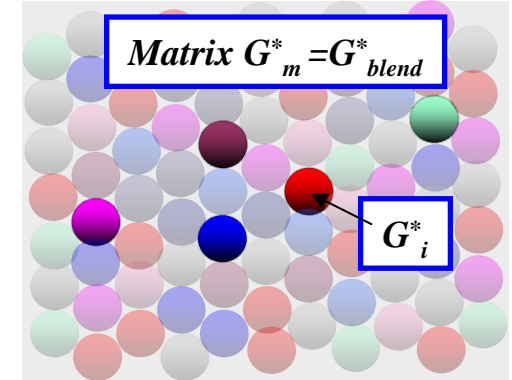
$$G_{blend}^* = G_m^* \left( 1 + \frac{5}{2} \sum_i f_i H_i \right)$$

$$H_i = \frac{2G_i^* - 2G_m^*}{2G_i^* + 3G_m^*}$$

Volume fraction

Olroyd-Palierne model.

[Palierne, 1990]



Heterogeneous visco-elastic system:

$$\sum_i f_i H_i = 0$$

$$H_i = \frac{2G_i^* - 2G_m^*}{2G_i^* + 3G_m^*}$$

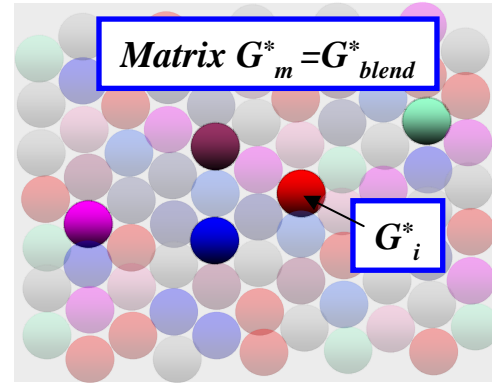
Self-consistent Olroyd-Palierne model.

[Lequeux, 2001]

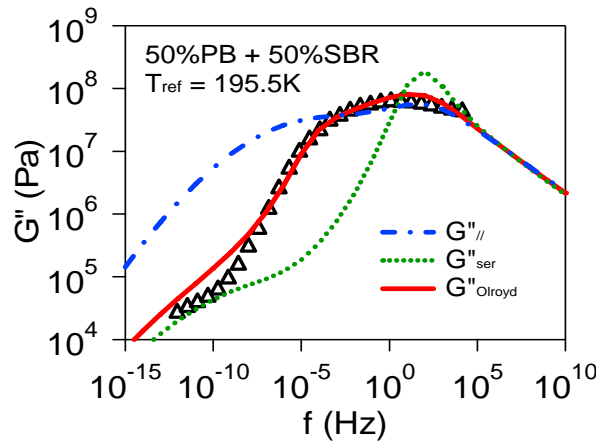
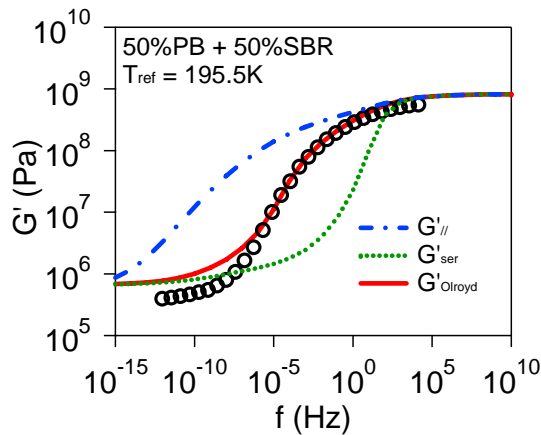
# Self-consistent Olroyd-Palierne model

- Self-consistent Olroyd-Palierne model : [Lequeux, 2001]

$$\int \frac{2G^*(T_g) - 2G_{blend}^*}{2G^*(T_g) + 3G_{blend}^*} P(T_g) dT_g = 0$$

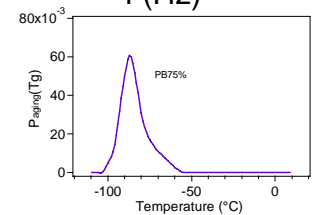
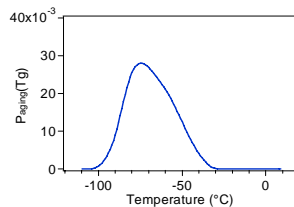
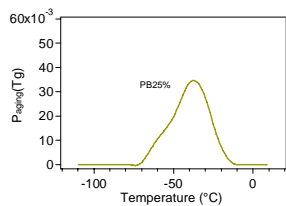
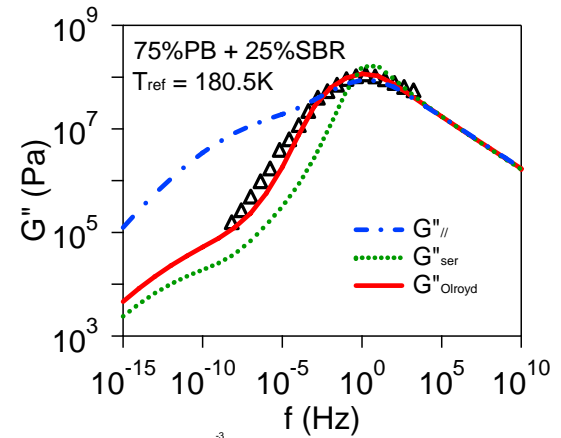
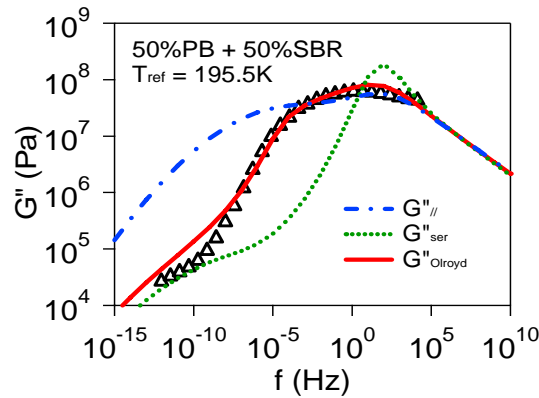
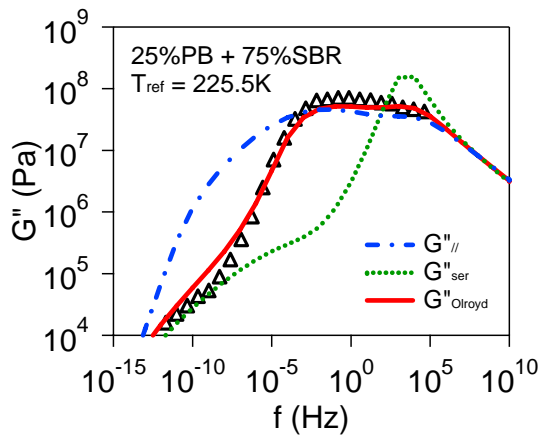
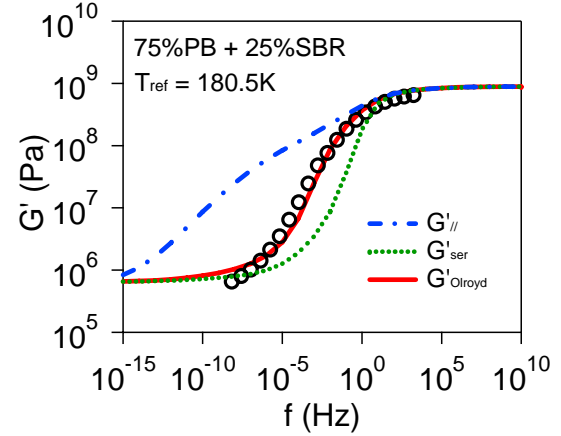
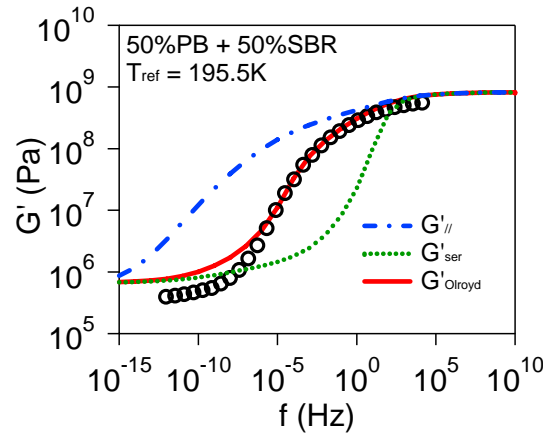
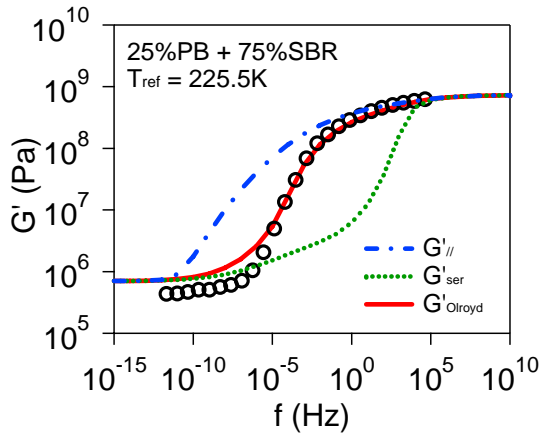


*Spherical inclusions  $G_i$   
with volume fraction  $f_i$*

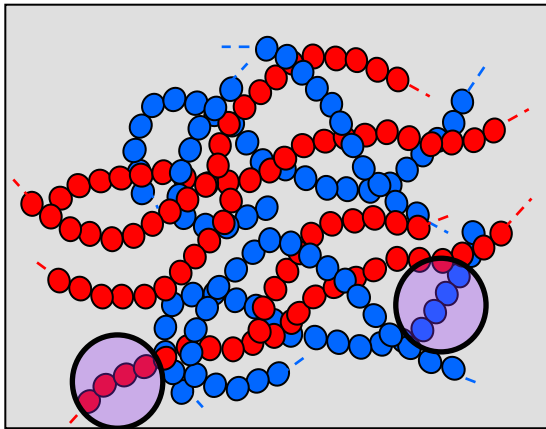


- Agreement!!
- Prediction of blend's rheology from DSC measurements and pure polymers' rheology.
- No adjustable parameters.

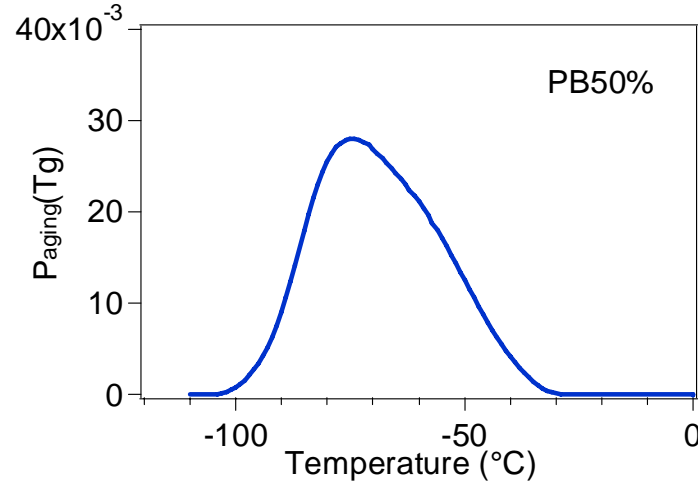
# Self-consistent Olroyd-Palierne model



# Rheology prediction



DSC  
Aging



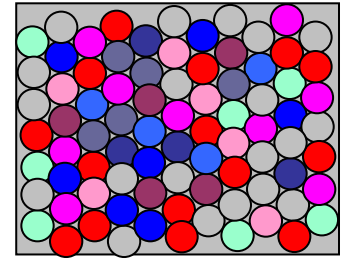
Rheology:OK!

## Conclusion 3:

- The rheology property of our blend is controlled by the same size scale (1nm), and the same Tg distribution.
- Our blend system can be considered in its glass transition regime as an arrangement of independent viscoelastic domains, each one with a specific glass transition. The best way to average viscoelastic moduli in such heterogeneous system, is the self-consistent Olroyd-Palierne model.



# General Conclusions



- **Physical aging** of our blend system occurs independently at domains of a scale  $\sim$  1nm.
- The **rheology** property of our blend system in its glass transition regime can be regarded as an ensemble of independent viscoelastic domains of a scale  $\sim$  1nm.
- Open question: is it general for polymer blends?

Accepted by *Macromolecules*

# Thanks:



- François LEQUEUX
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- Etienne MUNCH
- Régis SCHACH

Thank you!

